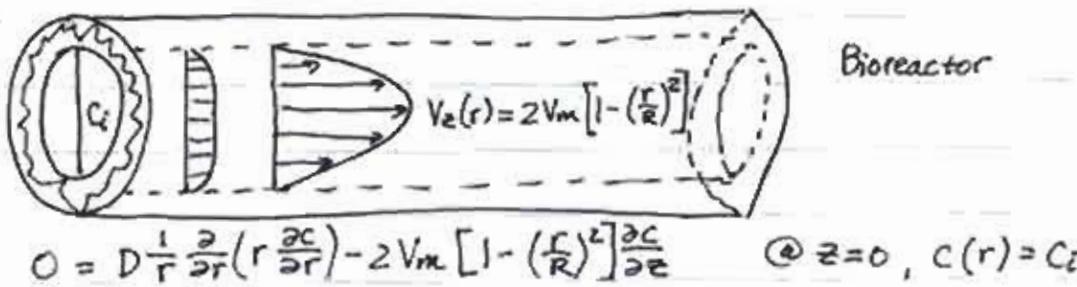


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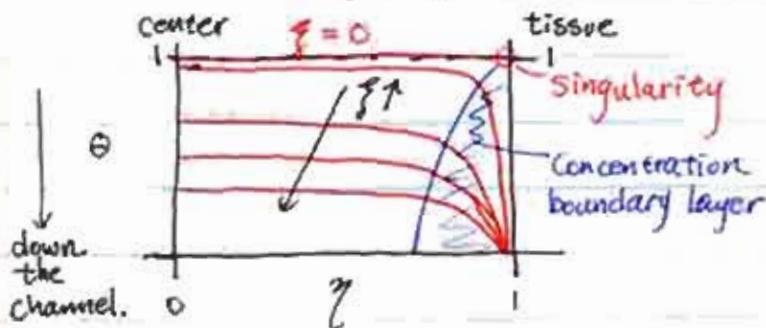
$$r=0, \frac{\partial c}{\partial r}$$

$$r=R, c=0.$$

Solution: $c(r, z) = c_i \Theta \left(\frac{r}{R}, \frac{z}{L} \right)$

$$= \sum_{j=1}^{\infty} \left[\int_0^1 \Phi_j(\gamma) (1-\gamma^2) \gamma d\gamma \right] \Phi_j(\eta) e^{-\frac{1}{Pe} \lambda_j^2 \frac{z}{L}}$$

$\Phi_j(\gamma), \lambda_j$ determined from corresponding E.V. Egn.



Objective: design to ensure adequate flux of nutrient into tissue throughout channel

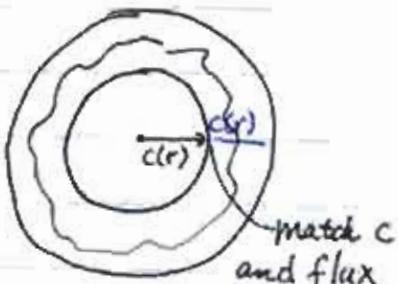
Flux @ wall $Nr/R(\xi) = -D \frac{\partial c}{\partial r} / R$ proportional to slope in boundary layer.
moles/area-time can calculate from $C(r, z)$

Might desire thickness of tissue region

$$Nr/R(\xi) > n \delta \mu_R$$

↑ cells/vol metabolic uptake rate
in tissue region moles/cell-time

Cross-section

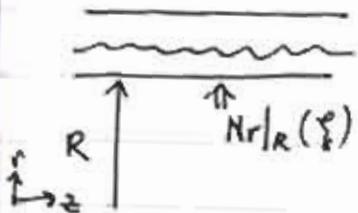
match c
and flux

Offer an alternative perspective

- "mass transfer coefficient" concept.

avg. conc.
in bulk fluid

[moles/ml]



Define: $Nr/R(z) = k_c (C_{bulk} - C_{wall}) \quad \text{in our problem}$

mass transfer coefficient.
 Conc. @ wall
 [distance/time] [moles/ml]

$\int_0^R V_z(r) \frac{\partial c}{\partial r} dr = \int_0^R D \frac{\partial}{\partial r} \left(r \frac{\partial c}{\partial r} \right) dr$ Take integral of both sides over a cylindrical shell in channel.



$$\text{Obtain LHS } \frac{d}{dz} \int_0^R V_z(r) + C(r) dr = \frac{1}{2} R^2 V_m \frac{dC_{bulk}}{dz}$$

$$C_{bulk} = \frac{\int_0^R C(r) V_z(r) r dr}{\int_0^R V_z(r) r dr} \quad \text{"velocity weighted" - bulk concentration}$$

$$\text{RHS: } \int_0^R D \frac{\partial}{\partial r} \left(r \frac{\partial c}{\partial r} \right) dr = D r \frac{\partial c}{\partial r} \Big|_0^R = DR \frac{\partial c}{\partial r} \Big|_R = -RN_r \Big|_R (z)$$

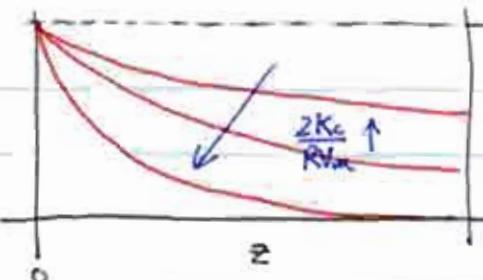
$$\text{So, } \frac{1}{2} R^2 V_m \frac{dC_{bulk}}{dz} = -R [k_c (C_B - C_w)] \quad // \text{Assume } k_c \neq k_c(z)$$

$$\Rightarrow \frac{dC_{bulk}}{dz} = \frac{-2k_c}{RV_m} (C_B - C_w) \quad z=0, C_{bulk}=C_i$$

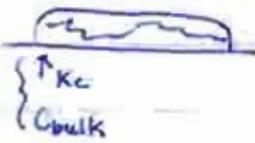
Solution: $\frac{C_{bulk} - C_w}{C_i - C_w} = \exp \left\{ -\frac{2k_c}{RV_m} z \right\} \quad \text{if } C_w \neq 0$

$\frac{-2k_c}{RV_m} z$

if $C_w = 0$ then $C_{bulk}(z) = C_i e^{-\frac{2k_c}{RV_m} z}$



But, what is K_c ?



$\frac{\text{distance}}{\text{time}}$

\downarrow $\frac{\text{distance}^2}{\text{time}}$

$\frac{\text{distance}^2}{\text{distance}}$

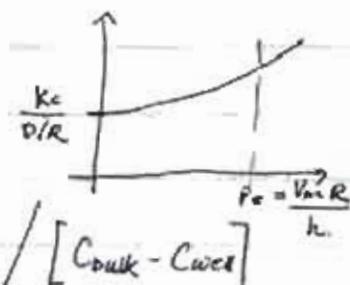
$$Nr = \frac{D \alpha c}{R}$$

If $V_m = 0$, have only diffusion, then $K_c = \frac{D}{R}$; $Nr = \frac{D \alpha c}{R}$



How to ascertain K_c under any condition.

Formally, Sherwood number : $Sh = \frac{K_c}{D/R} = -\frac{\partial c}{\partial r} \Big|_{\text{wall}} / [C_{\text{bulk}} - C_{\text{wall}}]$



How to determine Sh ?

a) Theory

b) Computation - model (yield a number from numerical solution)

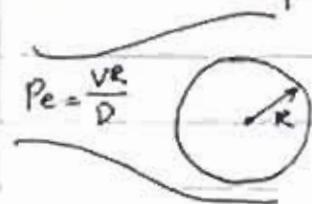
c) Experimental measurements - correlations w/ system parameters.

See sections 9.3 in Deen's text

Examples of a) Theory - problem that we did on Monday (Gratz problem)

$$\Rightarrow Sh = \frac{K_c}{D/R} \approx \frac{1}{2} \lambda_1^2 = 3.7 \leftarrow \text{smallest eigenvalue: as } z \rightarrow \infty$$

Flow around sphere



$$Sh = \frac{K_c}{V/2R} \quad Pe \ll 1$$