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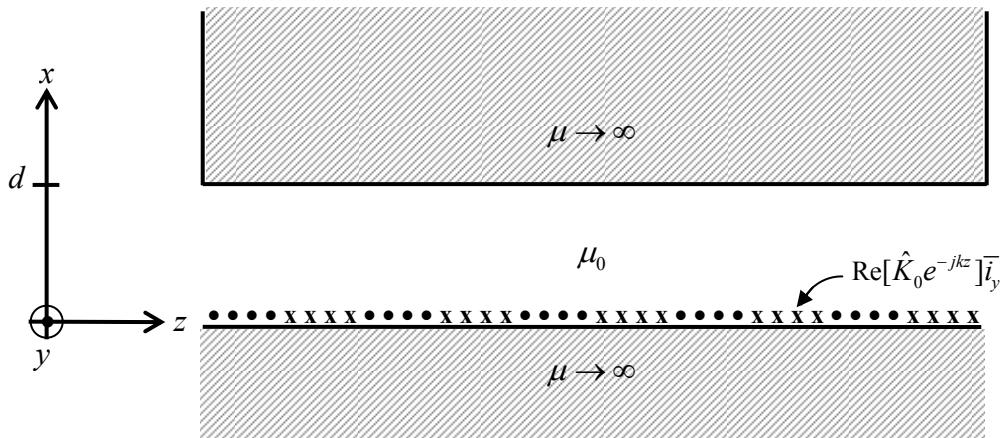
6.642 Continuum Electromechanics
Fall 2008

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Continuum Electromechanics
6.642 Mid-Term Exam
November 5, 2008

Formula sheets are located after page 4.

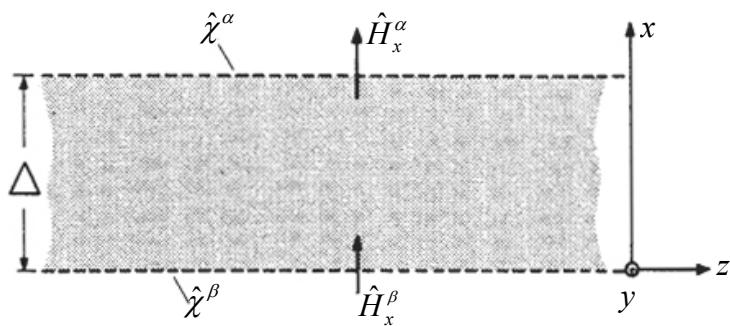
1. (30 points)



A current sheet $\text{Re}[\hat{K}_0 e^{-jkz}] \bar{i}_y$ is placed on the $x = 0$ surface of a material with infinite magnetic permeability ($\mu \rightarrow \infty$) for $x < 0$. Another infinite magnetic permeability material extends from $d < x < \infty$.

Free space with magnetic permeability μ_0 extends over the region $0 < x < d$.

The magnetic field (H_x) – magnetic scalar potential (χ) relations for the planar layer below



for variables of the form

$$\chi(x, z) = \text{Re}[\hat{\chi}(x) e^{-jkz}]$$

are

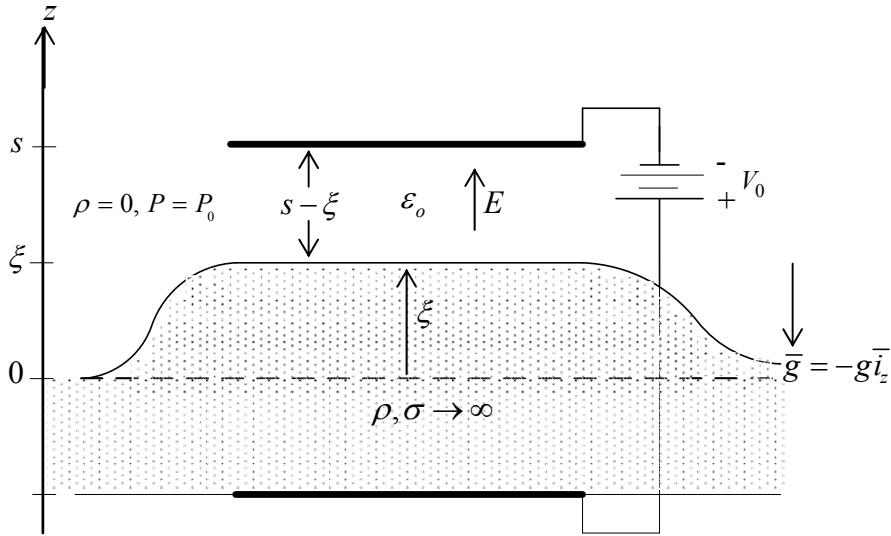
$$\begin{bmatrix} \hat{H}_x^\alpha \\ \hat{H}_x^\beta \end{bmatrix} = k \begin{bmatrix} -\coth k\Delta & \frac{1}{\sinh k\Delta} \\ -\frac{1}{\sinh k\Delta} & \coth k\Delta \end{bmatrix} \begin{bmatrix} \hat{\chi}^\alpha \\ \hat{\chi}^\beta \end{bmatrix}$$

where $\bar{H}(x, z) = -\nabla \chi(x, z) = \text{Re}[(\hat{H}_x(x)\bar{i}_x + \hat{H}_z(x)\bar{i}_z)e^{-jkz}]$

There is no magnetic field dependence on y.

- What are the boundary conditions on the magnetic field at the $x = 0_+$ and $x = d_-$ surfaces? What are the values of the magnetic scalar potential $\hat{\chi}(x = 0_+)$ and $\hat{\chi}(x = d_-)$?
- What are the complex amplitudes of the magnetic field $\bar{H}(x, z)$ at $x = 0_+$ and at $x = d_-$?
- What is the magnetic force per unit area (on a wave length $2\pi/k$) \bar{F} on the infinite magnetic permeability layer that extends $d < x < \infty$?

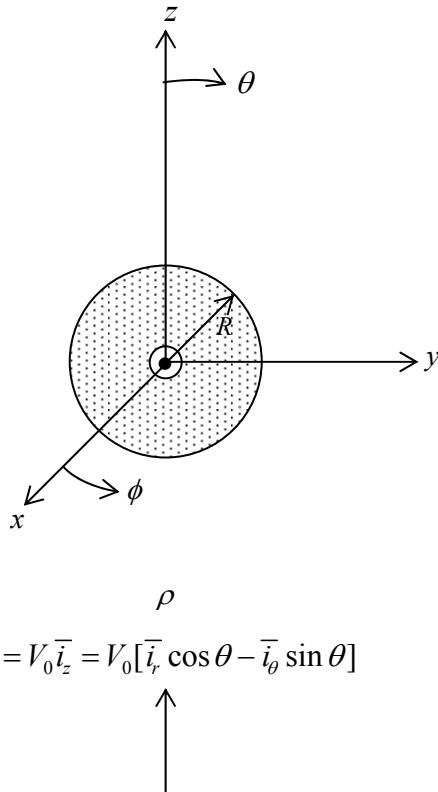
2. (35 points)



A perfectly conducting incompressible liquid ($\sigma \rightarrow \infty$) with mass density ρ partially fills the gap between parallel plate electrodes stressed by voltage V_0 . The applied voltage lifts the fluid interface between the parallel plate electrodes by a height ξ where $\xi < s$. The upper electrode in free space is at $z = s$. When the applied voltage is zero the fluid interface is located at $z = 0$. The region outside the liquid is free space with permittivity ϵ_o , mass density of zero ($\rho = 0$), and atmospheric pressure P_0 . The gravitational acceleration is $\bar{g} = -g\bar{i}_z$ and surface tension effects are negligible.

- What is the electric field for $\xi < z < s$ between the upper electrode at $z = s$ and the perfectly conducting fluid interface at $z = \xi$?
- What is the fluid pressure $p(\xi)$ just below the interface at $z = \xi$?
- Find an expression that relates liquid rise ξ ($\xi < s$) to voltage V_0 and other given parameters.
- At what voltage is $\xi = \frac{s}{2}$?

3. (35 points)



An inviscid incompressible liquid with mass density ρ has uniform irrotational flow ($\nabla \times \bar{v} = 0$). The flow at $r = \infty$ is uniform and z directed

$$\bar{v} = V_0 \bar{i}_z = V_0 [\bar{i}_r \cos \theta - \bar{i}_\theta \sin \theta]$$

The flow is incident on a solid sphere of radius R . The inviscid liquid can flow along the sphere so that $v_\theta(r = R_+) \neq 0$ but cannot penetrate the surface so that $v_r(r = R_+) = 0$. Because the irrotational flow has $\nabla \times \bar{v} = 0$, a velocity scalar potential Φ can be defined, $\bar{v} = -\nabla \Phi$. Because the fluid is also incompressible, $\nabla \cdot \bar{v} = 0$, the velocity scalar potential for $r > R$ obeys Laplace's equation, $\nabla^2 \Phi = 0$ where $\Phi(r, \theta)$ does not depend on angle ϕ . The flow does not vary with time and gravity effects are negligible.

- a) What are the boundary conditions on the velocity scalar potential at $r = R_+$ and at $r = \infty$?
- b) Solve for the velocity scalar potential $\Phi(r, \theta)$.
- c) Solve for the velocity field $\bar{v}(r, \theta)$ for $r > R$.
- d) What is the magnitude of the velocity $|\bar{v}(r, \theta)|$?
- e) If the pressure at $r = R_+$ and $\theta = 0$ is P_0 , what is the pressure at $r = R_+$, $\theta = \pi/2$?
- f) What is the equation for the velocity streamlines?

Hint: $\int \frac{\cos d\theta}{\sin \theta} = \ln(\sin \theta) + \text{constant}$

$$\int \frac{1 + \frac{R^3}{2r^3}}{r(1 - \frac{R^3}{r^3})} dr = \frac{1}{2} \ln[r^2(1 - \frac{R^3}{r^3})] + \text{constant}$$

- g) For the velocity streamline that passes through the point $x = 0$, $y = y_0$, $z = 0$ equivalent to $r = y_0$, $\theta = \frac{\pi}{2}$, $\phi = \frac{\pi}{2}$, for what value of y does the streamline pass through when $x = 0$ and $z = -\infty$, equivalent to $r = \infty$, $\theta = \pi$, $\phi = \frac{\pi}{2}$? Find y when $y_0 = R$ and when $y_0 = 2R$.

Hint: $y = r \sin \theta \sin \phi$

6.642 Formula Sheet #1

Cartesian Coordinates (x, y, z)

$$\begin{aligned}\nabla f &= \frac{\partial f}{\partial x} \mathbf{i}_x + \frac{\partial f}{\partial y} \mathbf{i}_y + \frac{\partial f}{\partial z} \mathbf{i}_z \\ \nabla \cdot A &= \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z} \\ \nabla \times A &= \mathbf{i}_x \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) + \mathbf{i}_y \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) + \mathbf{i}_z \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \\ \nabla^2 f &= \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2}\end{aligned}$$

Cylindrical Coordinates (r, ϕ, z)

$$\begin{aligned}\nabla f &= \frac{\partial f}{\partial r} \mathbf{i}_r + \frac{1}{r} \frac{\partial f}{\partial \phi} \mathbf{i}_\phi + \frac{\partial f}{\partial z} \mathbf{i}_z \\ \nabla \cdot A &= \frac{1}{r} \frac{\partial}{\partial r} (r A_r) + \frac{1}{r} \frac{\partial A_\phi}{\partial \phi} + \frac{\partial A_z}{\partial z} \\ \nabla \times A &= \mathbf{i}_r \left(\frac{1}{r} \frac{\partial A_z}{\partial \phi} - \frac{\partial A_\phi}{\partial z} \right) + \mathbf{i}_\phi \left(\frac{\partial A_r}{\partial z} - \frac{\partial A_z}{\partial r} \right) + \mathbf{i}_z \frac{1}{r} \left[\frac{\partial (r A_\phi)}{\partial r} - \frac{\partial A_r}{\partial \phi} \right] \\ \nabla^2 f &= \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial f}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 f}{\partial \phi^2} + \frac{\partial^2 f}{\partial z^2}\end{aligned}$$

Spherical Coordinates (r, θ, ϕ)

$$\begin{aligned}\nabla f &= \frac{\partial f}{\partial r} \mathbf{i}_r + \frac{1}{r} \frac{\partial f}{\partial \theta} \mathbf{i}_\theta + \frac{1}{r \sin \theta} \frac{\partial f}{\partial \phi} \mathbf{i}_\phi \\ \nabla \cdot A &= \frac{1}{r^2} \frac{\partial}{\partial (r A_r)} + \frac{1}{r \sin \theta} \frac{\partial (r A_\theta)}{\partial \theta} + \frac{1}{\sin \theta} \frac{\partial A_\phi}{\partial \phi} \\ \nabla \times A &= \mathbf{i}_r \frac{1}{r} \left[\frac{\partial (\sin \theta A_\phi)}{\partial \theta} - \frac{\partial A_\theta}{\partial \phi} \right] \\ &\quad + \mathbf{i}_\theta \frac{1}{r} \left[\frac{1}{\sin \theta} \frac{\partial A_r}{\partial \phi} - \frac{\partial (r A_\phi)}{\partial r} \right] + \mathbf{i}_\phi \frac{1}{r} \left[\frac{\partial (r A_\theta)}{\partial r} - \frac{\partial A_r}{\partial \theta} \right] \\ \nabla^2 f &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial f}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial f}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 f}{\partial \phi^2}\end{aligned}$$

Cartesian

Cylindrical

Spherical

$$\begin{aligned}x &= r \cos \phi & r \cos \phi &= r \sin \theta \cos \phi \\ y &= r \sin \phi & r \sin \phi &= r \sin \theta \sin \phi \\ z &= z & z &= r \cos \theta \\ \mathbf{i}_x &= \cos \phi \mathbf{i}_r - \sin \phi \mathbf{i}_\phi & \cos \phi \mathbf{i}_r &= \sin \theta \cos \phi \mathbf{i}_r + \cos \theta \cos \phi \mathbf{i}_\theta \\ \mathbf{i}_y &= \sin \phi \mathbf{i}_r + \cos \phi \mathbf{i}_\phi & \sin \phi \mathbf{i}_r &= \sin \theta \sin \phi \mathbf{i}_r + \cos \theta \sin \phi \mathbf{i}_\theta \\ \mathbf{i}_z &= \mathbf{i}_z & \mathbf{i}_z &= \cos \theta \mathbf{i}_r - \sin \theta \mathbf{i}_\theta\end{aligned}$$

Cylindrical

Cartesian

Spherical

$$\begin{aligned}r &= \sqrt{x^2 + y^2 + z^2} & r \cos \phi &= \sqrt{r^2 + z^2} \\ \theta &= \cos^{-1} \frac{z}{\sqrt{x^2 + y^2 + z^2}} & r \sin \phi &= \cos^{-1} \frac{z}{\sqrt{r^2 + z^2}} \\ \phi &= \cot^{-1} \frac{y}{x} & \phi &= \cot^{-1} \frac{y}{\sqrt{r^2 + z^2}} \\ \mathbf{i}_r &= \sin \theta \cos \phi \mathbf{i}_x + \sin \theta \sin \phi \mathbf{i}_y + \cos \theta \mathbf{i}_z & \sin \theta \mathbf{i}_r + \cos \theta \mathbf{i}_z &= \sin \theta \mathbf{i}_r + \cos \theta \mathbf{i}_z \\ \mathbf{i}_\theta &= \cos \theta \cos \phi \mathbf{i}_x + \cos \theta \sin \phi \mathbf{i}_y - \sin \theta \mathbf{i}_z & \cos \theta \mathbf{i}_r - \sin \theta \mathbf{i}_z &= \cos \theta \mathbf{i}_r - \sin \theta \mathbf{i}_z\end{aligned}$$

Geometric relations between coordinates and unit vectors for Cartesian, cylindrical, and spherical coordinate systems.

6.642 Formula Sheet #1

Vector Identities	Maxwell's Equations	Boundary Conditions
<i>Integral</i>	<i>Differential</i>	
$(\mathbf{A} \times \mathbf{B}) \cdot \mathbf{C} = \mathbf{A} \cdot (\mathbf{B} \times \mathbf{C}) = (\mathbf{C} \times \mathbf{A}) \cdot \mathbf{B}$ $\mathbf{A} \times (\mathbf{B} \times \mathbf{C}) = \mathbf{B}(\mathbf{A} \cdot \mathbf{C}) - \mathbf{C}(\mathbf{A} \cdot \mathbf{B})$ $\nabla \cdot (\nabla \times \mathbf{A}) = 0$ $\nabla \times (\nabla f) = 0$ $\nabla(fg) = f\nabla g + g\nabla f$ $\nabla(\mathbf{A} \cdot \mathbf{B}) = (\mathbf{A} \cdot \nabla)\mathbf{B} + (\mathbf{B} \cdot \nabla)\mathbf{A} + \mathbf{A} \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{A})$ $\nabla \cdot (f\mathbf{A}) = f\nabla \cdot \mathbf{A} + (\mathbf{A} \cdot \nabla)f$ $\nabla \cdot (\mathbf{A} \times \mathbf{B}) = \mathbf{B} \cdot (\nabla \times \mathbf{A}) - \mathbf{A} \cdot (\nabla \times \mathbf{B})$ $\nabla \times (\mathbf{A} \times \mathbf{B}) = \mathbf{A}(\nabla \cdot \mathbf{B}) - \mathbf{B}(\nabla \cdot \mathbf{A}) + (\mathbf{B} \cdot \nabla)\mathbf{A} - (\mathbf{A} \cdot \nabla)\mathbf{B}$ $\nabla \times (f\mathbf{A}) = \nabla f \times \mathbf{A} + f\nabla \times \mathbf{A}$ $(\nabla \times \mathbf{A}) \times \mathbf{A} = (\mathbf{A} \cdot \nabla)\mathbf{A} - \frac{1}{2}\nabla(\mathbf{A} \cdot \mathbf{A})$ $\nabla \times (\nabla \times \mathbf{A}) = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$	$\int_L \mathbf{E}' \cdot d\mathbf{l} = -\frac{d}{dt} \int_S \mathbf{B} \cdot d\mathbf{s}$ Faraday's Law $\int_L \mathbf{H} \cdot d\mathbf{l} = \int_S \mathbf{J}_f \cdot d\mathbf{s} + \frac{d}{dt} \int_S \mathbf{D} \cdot d\mathbf{s}$ Ampere's Law with Maxwell's Displacement Current Correction $\int_S \mathbf{D} \cdot d\mathbf{s} = \int_V \rho_f dV$ Gauss's Law $\int_S \mathbf{B} \cdot d\mathbf{s} = 0$ Conservation of Charge $\int_S \mathbf{J}_f \cdot d\mathbf{s} + \frac{d}{dt} \int_V \rho_f dV' = 0$ Usual Linear Constitutive Laws $\mathbf{D} = \epsilon \mathbf{E}$ $\mathbf{B} = \mu \mathbf{H}$ $\mathbf{J}_f = \sigma(\mathbf{E} + \mathbf{v} \times \mathbf{B}) = \sigma \mathbf{E}'$ [Ohm's law for moving media with velocity \mathbf{v}]	$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ $\mathbf{n} \times (\mathbf{E}_2' - \mathbf{E}_1') = 0$ $\nabla \times \mathbf{H} = \mathbf{J}_f + \frac{\partial \mathbf{D}}{\partial t}$ $\mathbf{n} \times (\mathbf{H}_2 - \mathbf{H}_1) = \mathbf{K}_f$ $\nabla \cdot \mathbf{D} = \rho_f$ $\nabla \cdot \mathbf{B} = 0$ $\mathbf{n} \cdot (\mathbf{D}_2 - \mathbf{D}_1) = \sigma_f$ $\mathbf{n} \cdot (\mathbf{B}_2 - \mathbf{B}_1) = 0$ $\nabla \cdot \mathbf{J}_f + \frac{\partial \rho_f}{\partial t} = 0$ $\mathbf{n} \cdot (\mathbf{J}_2 - \mathbf{J}_1) + \frac{\partial \sigma_f}{\partial t} = 0$
Integral Theorems		
<i>Line Integral of a Gradient</i> $\int_a^b \nabla f \cdot d\mathbf{l} = f(b) - f(a)$		
<i>Divergence Theorem:</i> $\int_V \nabla \cdot \mathbf{A} dV = \oint_S \mathbf{A} \cdot d\mathbf{s}$		
<i>Corollaries</i>		
$\int_V \nabla f dV = \oint_S f d\mathbf{s}$ $\int_V \nabla \times \mathbf{A} dV = - \oint_S \mathbf{A} \times d\mathbf{s}$		
<i>Stokes' Theorem:</i> $\oint_L \mathbf{A} \cdot d\mathbf{l} = \int_S (\nabla \times \mathbf{A}) \cdot d\mathbf{s}$		
<i>Corollary</i> $\oint_L f d\mathbf{l} = - \int_S \nabla f \times d\mathbf{s}$		

APPENDIX A. Differential Operators in Cartesian, Cylindrical and Spherical Coordinates

Operator	Cartesian coordinates	Cylindrical coordinates	Spherical coordinates
$(\nabla \cdot \vec{A})$	$\frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$	$\frac{1}{r} \frac{\partial}{\partial r} (r A_r) + \frac{1}{r} \frac{\partial A_\theta}{\partial \theta} + \frac{\partial A_\phi}{\partial z}$	$\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (A_\theta \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial A_\phi}{\partial \phi}$
$\nabla \Phi$	$\frac{\partial \Phi}{\partial x} \hat{i}_x + \frac{\partial \Phi}{\partial y} \hat{i}_y + \frac{\partial \Phi}{\partial z} \hat{i}_z$	$\frac{\partial \Phi}{\partial r} \hat{i}_r + \frac{1}{r} \frac{\partial \Phi}{\partial \theta} \hat{i}_\theta + \frac{\partial \Phi}{\partial z} \hat{i}_z$	$\frac{\partial \Phi}{\partial r} \hat{i}_r + \frac{1}{r} \frac{\partial \Phi}{\partial \theta} \hat{i}_\theta + \frac{1}{r \sin \theta} \frac{\partial \Phi}{\partial \phi} \hat{i}_\phi$
$(\nabla^2 \Phi)$	$\frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial y^2} + \frac{\partial^2 \Phi}{\partial z^2}$	$\frac{1}{r} \frac{\partial}{\partial r} (r \frac{\partial \Phi}{\partial r}) + \frac{1}{r^2} \frac{\partial^2 \Phi}{\partial \theta^2} + \frac{\partial^2 \Phi}{\partial z^2}$	$\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \frac{\partial \Phi}{\partial r}) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta \frac{\partial \Phi}{\partial \theta}) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \Phi}{\partial \phi^2}$
$(\nabla \times \vec{A})$	$\left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) \hat{i}_x + \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) \hat{i}_y + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \hat{i}_z$	$\left(\frac{1}{r} \frac{\partial A_\phi}{\partial \theta} - \frac{\partial A_\theta}{\partial z} \right) \hat{i}_r + \left(\frac{\partial A_r}{\partial z} - \frac{\partial A_z}{\partial r} \right) \hat{i}_\theta + \left(\frac{1}{r} \frac{\partial}{\partial r} (r A_\theta) - \frac{1}{r} \frac{\partial A_r}{\partial \theta} \right) \hat{i}_\phi$	$\left(\frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (A_\phi \sin \theta) - \frac{1}{r \sin \theta} \frac{\partial A_\theta}{\partial \phi} \right) \hat{i}_r + \left(\frac{1}{r \sin \theta} \frac{\partial A_r}{\partial \phi} - \frac{1}{r} \frac{\partial}{\partial r} (r A_\phi) \right) \hat{i}_\theta + \left(\frac{1}{r} \frac{\partial}{\partial r} (r A_\theta) - \frac{1}{r} \frac{\partial A_r}{\partial \theta} \right) \hat{i}_\phi$

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Appendix A in Melcher, James R. *Continuum Electromechanics*. Cambridge, MA: MIT Press, 1981. ISBN: 9780262131650.

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Massachusetts Institute of Technology
 Department of Electrical Engineering and Computer Science
6.642 FORMULA SHEET

1. DIFFERENTIAL OPERATORS IN CYLINDRICAL AND SPHERICAL COORDINATES

If r , ϕ , and z are circular [cylindrical coordinates] and \hat{i}_r , \hat{i}_ϕ , and \hat{i}_z are unit vectors in the directions of increasing values of the corresponding coordinates,

$$\nabla U = \text{grad } U = \hat{i}_r \frac{\partial U}{\partial r} + \hat{i}_\phi \frac{1}{r} \frac{\partial U}{\partial \phi} + \hat{i}_z \frac{\partial U}{\partial z}$$

$$\nabla \cdot \vec{A} = \text{div } \vec{A} = \frac{1}{r} \frac{\partial (r A_r)}{\partial r} + \frac{1}{r} \frac{\partial A_\phi}{\partial \phi} + \frac{\partial A_z}{\partial z}$$

$$\nabla \times \vec{A} = \text{curl } \vec{A} = \hat{i}_r \left(\frac{1}{r} \frac{\partial A_z}{\partial \phi} - \frac{\partial A_\phi}{\partial z} \right) + \hat{i}_\phi \left(\frac{\partial A_r}{\partial z} - \frac{\partial A_z}{\partial r} \right) + \hat{i}_z \left(\frac{1}{r} \frac{\partial (r A_\phi)}{\partial r} - \frac{1}{r} \frac{\partial A_r}{\partial \phi} \right)$$

$$\nabla^2 U = \text{div grad } U = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial U}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 U}{\partial \phi^2} + \frac{\partial^2 U}{\partial z^2}$$

If r , θ , and ϕ are [spherical coordinates] and \hat{i}_r , \hat{i}_θ , and \hat{i}_ϕ are unit vectors in the directions of increasing values of the corresponding coordinates,

$$\nabla U = \text{grad } U = \hat{i}_r \frac{\partial U}{\partial r} + \hat{i}_\theta \frac{1}{r} \frac{\partial U}{\partial \theta} + \hat{i}_\phi \frac{1}{r \sin \theta} \frac{\partial U}{\partial \phi}$$

$$\nabla \cdot \vec{A} = \text{div } \vec{A} = \frac{1}{r^2} \frac{\partial (r^2 A_r)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial (A_\theta \sin \theta)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial A_\phi}{\partial \phi}$$

$$\nabla \times \vec{A} = \text{curl } \vec{A} = \hat{i}_r \left(\frac{1}{r \sin \theta} \frac{\partial (A_\phi \sin \theta)}{\partial \theta} - \frac{1}{r \sin \theta} \frac{\partial A_\theta}{\partial \phi} \right) + \hat{i}_\theta \left(\frac{1}{r \sin \theta} \frac{\partial A_r}{\partial \phi} - \frac{1}{r} \frac{\partial (r A_\phi)}{\partial r} \right) + \hat{i}_\phi \left(\frac{1}{r} \frac{\partial (r A_\theta)}{\partial r} - \frac{1}{r} \frac{\partial A_r}{\partial \theta} \right)$$

$$\nabla^2 U = \text{div grad } U = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial U}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial U}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 U}{\partial \phi^2}$$

2. SOLUTIONS OF LAPLACE'S EQUATIONS

A. Rectangular coordinates, two dimensions (independent of z):

$$\Phi = e^{kx} (A_1 \sin ky + A_2 \cos ky) + e^{-kx} (B_1 \sin ky + B_2 \cos ky)$$

(or replace e^{kx} and e^{-kx} by $\sinh kx$ and $\cosh kx$).

$$\Phi = Axy + Bx + Cy + D; (k = 0)$$

B. Cylindrical coordinates, two dimensions (independent of z):

$$\Phi = r^n (A_1 \sin n\phi + A_2 \cos n\phi) + r^{-n} (B_1 \sin n\phi + B_2 \cos n\phi)$$

$$\Phi = \ln \frac{R}{r} (A_1 \phi + A_2) + B_1 \phi + B_2; (n = 0)$$

C. Spherical coordinates, two dimensions (independent of ϕ):

$$\Phi = Ar \cos \theta + \frac{B}{r^2} \cos \theta + \frac{C}{r} + D$$